

## Relativistic Quantum Mechanics :-

### Klein - Gordon

It is the combination of (Quantum mechanics + special theory of relativity). This theory is highly accurate. In this theory the motion of that particle is considered whose speed is comparable to the speed of light.

### Klein - Gordon eq<sup>n</sup> :-

This eq<sup>n</sup> is also known as relativistic wave eq<sup>n</sup> and it is equivalent to schroedinger's wave eq<sup>n</sup>. It is second order differentiation eq<sup>n</sup> in space & time, satisfying the Lorentz transformation. This eq<sup>n</sup> may be used for (non-relativistic) as well as fast (relativistic). It is only appreciable for spinless particle. The wave func<sup>n</sup> of K-G eq<sup>n</sup> is not same as obtained. This eq<sup>n</sup> does not contain potential term.

Schrödinger eqn is given by

$$\hat{H} \psi = J \hbar \frac{\partial \psi}{\partial t} \quad \textcircled{1}$$

$$\hat{H} = -\frac{\hat{p}^2}{2m} = -\frac{\hbar^2 \nabla^2}{2m} \quad \textcircled{2}$$

Substituting in ①, we get

$$-\frac{\hbar^2}{2m} \nabla^2 \psi(r_1, t) = J \hbar \frac{\partial \psi(r_1, t)}{\partial t} \quad \textcircled{3}$$

This eqn is invariant under Galilean transformations or Linear Transformations but it is not invariant under Lorentz transformation in Special theory of relativity.

Therefore, this eqn is not applicable in relativistic systems.

In relativity, the energy i.e; Hamiltonian 'H' for a free particle is given by:

$$\hat{H} = E = \pm \sqrt{(\hat{p}^2 c^2 + m^2 c^4)}$$

$$\hat{H} = E = \pm \left[ \hat{p}^2 c^2 + m^2 c^4 \right]^{\frac{1}{2}} \quad \textcircled{4}$$

The Schrödinger eq<sup>n</sup> becomes :-

$$\pm \left[ \hat{p}^2 c^2 + m^2 c^4 \right]^{1/2} \psi = i\hbar \frac{\partial \psi}{\partial t}$$

(5)

This is Schrödinger eq<sup>n</sup> in relativistic medium but this eq<sup>n</sup> could not satisfy the energy states as obtained by Bohr and Sommerfeld.

So, Schrödinger was disappointed and put aside this eq<sup>n</sup>.

After this, Klein & Gordon modified this eq<sup>n</sup>.

They found the following difficulties in this eq<sup>n</sup>:-(7)

1. The first difficulty is of '+ve' and '-ve' signs. In classical mechanics, this was not problem because +ve and -ve energies, both exists and therefore the -ve energies can be ignored. But in quantum mechanics '-ve' energy does not exist.

2. It is not possible to interpret the square root of an operator i.e., the operator can not exist in its square-root

form

3. This eqn is only applicable for spinless particle.

These above difficulties were removed approximately by Klein - Gordon. They operated the entire eqn by ' $H$ '.

From ①

$$\hat{A}\psi = i\hbar \frac{\partial \psi}{\partial t}$$

$$\hat{H}^2 \psi = \hat{A} \left( i\hbar \frac{\partial \psi}{\partial t} \right)$$

$$= i\hbar \frac{\partial}{\partial t} \left( \hat{A}\psi \right)$$

$$= i\hbar \frac{\partial}{\partial t} \left( i\hbar \frac{\partial \psi}{\partial t} \right)$$

$$\hat{H}^2 \psi = -\hbar^2 \frac{\partial^2 \psi}{\partial t^2}$$

⑥

where  $E^2 = H^2$  This is k.G. eqn  
for a free particle  $\hat{H}^2 = \hat{p}^2 c^2 + m^2 c^4$  or relativistic Schrödinger eqn

$$= \left( \frac{\hbar}{i} \nabla \right)^2 c^2 + m^2 c^4$$

$$H^2 = -\hbar^2 c^2 \nabla^2 + m^2 c^4 \quad (7)$$

put in (6)

$$(-\hbar^2 c^2 \nabla^2 + m^2 c^4) \psi = -\hbar^2 \frac{\partial^2 \psi}{\partial t^2}$$

$$\left[ \nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \frac{m^2 c^2}{\hbar^2} \right] \psi = 0 \quad (8)$$

This is Klein-Gordon eqn for a  
or free particle

Relativistic Schrödinger eqn  
or

Wave eqn for relativistic

This eqn contains  $\hbar$  (quantum mechanics) & systems.

and  $c$  (relativity) i.e; It reflects quantum mechanics as well as relativity.

It may be written as :-

$$\left[ \square^2 - \frac{m^2 c^2}{\hbar^2} \right] \psi = 0$$

where  $\square^2 = \nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2}$

$\square^2$  is known as D'Alembertian operator.

In co-ordinate representation equation (8) of K.G.

Equation is

$$g^{(1)} \frac{\partial}{\partial x^1} \cdot \frac{\partial}{\partial x^1} - k^2 \psi(x^1) = 0$$

Where  $k = \frac{mc}{\hbar}$ , in this form the invariance of K.G. comes in action.

6.

$$\nabla^2 - \frac{1}{c^2} \cdot \frac{\partial^2}{\partial t^2}$$

space                          ↑ time

$$\frac{\partial}{\partial x^2} - \frac{1}{c^2} \cdot \frac{\partial^2}{\partial t^2} \quad i.e.; x^2 = c^2 t^2$$

Same as Lorentz transformation.